

# Decay Characteristics of Neutron Excess Nickel Nuclei

Joseph J Bevelacqua

Bevelacqua Resources, 7531 Flint Crossing Circle SE, Owens Cross Roads, AL 35763 USA

\*Corresponding author: Joseph J. Bevelacqua, Bevelacqua Resources, 7531 Flint Crossing Circle SE, Owens Cross Roads, AL 35762 USA.

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## Abstract

The properties of neutron excess nickel nuclei are determined utilizing a single-particle model. Single particle model calculations predict that  $A = 81 - 95$  neutron excess nickel systems form bound systems that have limiting beta decay half-lives in the range of  $0.310 - 10.5$  ms. Model half-life results for the  $A = 81 - 86$  nickel nuclei are within a factor of two of the predictions of the Japanese Nuclear Data Compilation calculations. The single particle model calculations include alpha, beta, positron, electron capture, and spontaneous fission decay modes. Neutron emission decay modes that have short half-lives are not readily determined by the model, and were not evaluated. The omission of these short-lived neutron emission decay modes implies that the single particle model calculations could overestimate the half-lives of neutron excess  $A = 81 - 95$  nickel nuclei.

**Keywords:** Nucleosynthesis, Neutron Excess Nickel Nuclei, Beta Decay, Nuclear Structure.

## Introduction

Interest in neutron excess nuclei has stimulated both experimental and theoretical physics interest. Several physical processes generate neutron excess nuclei, but the r-process usually provides the most significant contribution. Production of neutron excess nuclei in mergers of astrophysical objects (e.g., black holes and neutron stars) is an active area of research in nuclear physics and astrophysics [1-30].

This paper continues the investigation of neutron excess nuclei by focusing on the  $Z = 28$  nickel systems. Neutron excess systems having  $Z = 28$ , and 30 were discussed in previous work [8-25, 29, 30]. Studies of these systems provide additional insight into nuclear systematics involving the various nucleosynthesis mechanisms and decay modes, and their associated variation with atomic and mass numbers.

## Computational Methodology

Methods for investigating neutron excess nuclei are provided in Refs. [8-30]. This paper follows the single particle methodology of Lukasiak and Sobiczewski and Petrovich et al. [27, 28]. Single particle energies of neutron excess nuclear systems are obtained by incorporating the numerical methods of Refs [31, 32].

The radial Schrödinger equation is utilized to determine binding energy of a nucleon interacting with a nuclear core [8-25, 29, 30].

$$[(\hbar^2/2\mu) (d^2/dr^2 - L(L+1)/r^2) - E_{\text{NLSJ}} - V_{\text{LSJ}}(r)] U_{\text{NLSJ}}(r) = 0 \quad (1)$$

In Eq. 1,  $E_{\text{NLSJ}}$  is the nucleon binding energy,  $r$  is the radial coordinate,  $V_{\text{LSJ}}(r)$  is the nuclear interaction, and  $U_{\text{NLSJ}}(r)$  is the radial wave function.  $L$ ,  $S$ , and  $J$  represent the orbital, spin, and total angular momentum quantum numbers, respectively. The model definition is completed by defining the radial quantum number ( $N$ ) and reduced mass ( $\mu$ ).

## Nuclear Interaction

The Rost interaction is selected for the nuclear interaction [33]. This interaction has a central strength

$$V_0 = 51.6 [1 \pm 0.73 (N - Z)/A] \text{ MeV} \quad (2)$$

In Eq. 2, the positive (negative) sign is assigned to protons (neutrons). The spin-orbit interaction strength ( $V_{\text{so}}$ ) is defined in terms of the central interaction strength and the multiplier  $\gamma$  [33]:

$$V_{\text{so}} = \gamma V_0 / 180 \quad (3)$$

Inclusion of the pairing correction interaction of Blomqvist and

Wahlborn completes the definition of the model interaction [34].

The difficulties in defining an appropriate nuclear interaction are outlined in Refs. 35 and 36. Ray and Hodgson and Schwierz, Wiedenhöver, and Volya note that modifications, unique to each nuclear system, are required to ensure an accurate representation of the experimental energy levels and decay characteristics [35, 36]. In view of the conclusions of Refs. 35 and 36 and the results of previous excess neutron system calculations the Rost central interaction strength ( $V_A$ ) is modified in the following manner [8-25, 29, 30].

$$V_A = V_0 \lambda [1 \pm a(A)] \text{ MeV} \quad (4)$$

Individual nuclear system characteristics are defined by incorporating a potential strength multiplier ( $\lambda$ ) and a factor  $[a(A)]$  to adjust the potential strength as a function of  $A$ . For nickel systems, the multiplier  $\lambda$  is selected to have the value of 1.5. This multiplier value is consistent with previous excess neutron nuclei calculations [8-25, 29, 30] that provided model results in agreement with available data [37-39].

### Model Limitations

Previous calculations provided a representative description of the various nuclear decay modes (e.g., alpha, beta, positron, electron capture, and spontaneous fission) that could be encountered in neutron excess nuclei [8-25, 29, 30, 40]. Neutron excess systems can also decay by neutron emission modes that

are not well-described by single-particle models. Since these neutron emission modes have very short half-lives, single-particle models will likely overestimate the lifetimes of neutron excess nuclei.

### Results and Discussion

Table 1 summarizes the complete set of  $95 \geq A \geq 69$  nickel isotopes considered in this paper. The  $95 \geq A \geq 69$  nickel nuclei occupy the  $1g_{9/2}$  ( $^{69}\text{Ni} - ^{78}\text{Ni}$ ),  $2d_{5/2}$  ( $^{79}\text{Ni} - ^{84}\text{Ni}$ ),  $1g_{7/2}$  ( $^{85}\text{Ni} - ^{92}\text{Ni}$ ),  $3s_{1/2}$  ( $^{93}\text{Ni} - ^{94}\text{Ni}$ ), and  $2d_{3/2}$  ( $^{95}\text{Ni}$ ) neutron single-particle levels. The heaviest observed nickel system is  $^{80}\text{Ni}$  [37-39]. In view of the paucity of experimental data, extrapolations of nuclear characteristics beyond  $A > 80$  become more uncertain.

### $80 \geq A \geq 69$ Nickel Isotopes with Experimental Half-Life Data

The limiting decay mode (i.e., the transition that has the shortest decay half-life) for  $69 \geq A \geq 80$  nickel isotopes observed experimentally is summarized in Table 1. For example, the  $^{73}\text{Ni}$  calculations include seven beta decay transitions (i.e., allowed  $1f_{7/2}(n)$  to  $1f_{5/2}(p)$  [5.57 min], allowed  $1f_{5/2}(n)$  to  $1f_{5/2}(p)$  [2.58 s], allowed  $2p_{3/2}(n)$  to  $2p_{3/2}(p)$  [2.13 s], allowed  $2p_{3/2}(n)$  to  $2p_{1/2}(p)$  [9.32 s], allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  [839 ms], allowed  $2p_{1/2}(n)$  to  $2p_{1/2}(p)$  [2.54 s], and first forbidden  $1f_{5/2}(n)$  to  $1g_{9/2}(p)$  [5.08 d]). For  $^{73}\text{Ni}$ , the allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  [839 ms] transition is the limiting beta decay mode.

**Table 1:** Calculated Single-Particle and Experimental Decay Properties of Nickel Nuclei with  $69 \leq A \leq 95$

Nuclide	a(A)	Half-Life (Decay Mode)	
		Experiment <sup>a,b,c</sup> /Theory <sup>d</sup>	This Work
$^{69}\text{Ni}$	-0.0350	11.4 s <sup>c</sup>	11.4 s ( $\beta^-$ ) <sup>e</sup>
$^{70}\text{Ni}$	-0.0371	6.0 s <sup>c</sup>	6.02 s ( $\beta^-$ ) <sup>e</sup>
$^{71}\text{Ni}$	-0.0350	2.56 s <sup>c</sup>	2.56 s ( $\beta^-$ ) <sup>e</sup>
$^{72}\text{Ni}$	-0.0360	1.57 s <sup>c</sup>	1.57 s ( $\beta^-$ ) <sup>e</sup>
$^{73}\text{Ni}$	-0.0333	840 ms <sup>c</sup>	839 ms ( $\beta^-$ ) <sup>e</sup>
$^{74}\text{Ni}$	-0.0310	507.7 ms <sup>c</sup>	507 ms ( $\beta^-$ ) <sup>e</sup>
$^{75}\text{Ni}$	-0.0294	331.6 ms <sup>c</sup>	331 ms ( $\beta^-$ ) <sup>e</sup>
$^{76}\text{Ni}$	-0.0285	234.6 ms <sup>c</sup>	234 ms ( $\beta^-$ ) <sup>e</sup>
$^{77}\text{Ni}$	-0.0254	158.9 ms <sup>c</sup>	159 ms ( $\beta^-$ ) <sup>e</sup>
$^{78}\text{Ni}$	-0.0249	122 ms <sup>c</sup>	122 ms ( $\beta^-$ ) <sup>e</sup>
$^{79}\text{Ni}$	+0.0047	43.0 ms <sup>c</sup>	43.0 ms ( $\beta^-$ ) <sup>e</sup>
$^{80}\text{Ni}$	+0.0232	23.9 ms <sup>c</sup>	23.9 ms ( $\beta^-$ ) <sup>e</sup>
$^{81}\text{Ni}$	+0.0417	12.3 ms <sup>d</sup>	10.5 ms ( $\beta^-$ ) <sup>e</sup>
$^{82}\text{Ni}$	+0.0602	8.74 ms <sup>d</sup>	9.17 ms ( $\beta^-$ ) <sup>e</sup>
$^{83}\text{Ni}$	+0.0787	7.57 ms <sup>d</sup>	6.12 ms ( $\beta^-$ ) <sup>e</sup>
$^{84}\text{Ni}$	+0.0972	5.56 ms <sup>d</sup>	4.25 ms ( $\beta^-$ ) <sup>e</sup>
$^{85}\text{Ni}$	+0.1157	4.74 ms <sup>d</sup>	3.06 ms ( $\beta^-$ ) <sup>e</sup>
$^{86}\text{Ni}$	+0.1342	3.45 ms <sup>d</sup>	2.26ms ( $\beta^-$ ) <sup>e</sup>
$^{87}\text{Ni}$	+0.1527	f	1.70 ms ( $\beta^-$ ) <sup>e</sup>
$^{88}\text{Ni}$	+0.1712	f	1.31 ms ( $\beta^-$ ) <sup>e</sup>
$^{89}\text{Ni}$	+0.1897	f	1.03 ms ( $\beta^-$ ) <sup>e</sup>
$^{90}\text{Ni}$	+0.2082	f	0.817 ms ( $\beta^-$ ) <sup>e</sup>

<sup>91</sup> Ni	+0.2267	f	0.659 ms ( $\beta^-$ ) <sup>c</sup>
<sup>92</sup> Ni	+0.2452	f	0.538 ms ( $\beta^-$ ) <sup>c</sup>
<sup>93</sup> Ni	+0.2637	f	0.444 ms ( $\beta^-$ ) <sup>c</sup>
<sup>94</sup> Ni	+0.2822	f	0.369 ms ( $\beta^-$ ) <sup>c</sup>
<sup>95</sup> Ni	+0.3007	f	0.310 ms ( $\beta^-$ ) <sup>c</sup>

<sup>a</sup>Ref. 37. <sup>b</sup>Ref. 38. <sup>c</sup>Ref. 39.

<sup>d</sup>Japanese data Compilation calculation (Ref. 39).

<sup>e</sup>Allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transition.

<sup>f</sup>No data provided in Ref. 37 - 39.

The model predicts the proper decay mode for the known  $80 \geq A \geq 69$  nickel systems [37 – 39]. As noted in Table 1, the model half-lives are also consistent with data [37 – 39].

<sup>69</sup>Ni – <sup>78</sup>Ni nuclei occupy the  $1g_{9/2}$  neutron shell. These systems decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta transitions. Model predictions for the beta decay half-lives of <sup>69</sup>Ni – <sup>78</sup>Ni are within 0.4% of the experimental values [39]. In addition beta decay is the predicted decay mode in agreement with Ref. 39.

The <sup>79</sup>Ni – <sup>80</sup>Ni nuclei partially fill the  $2d_{5/2}$  neutron shell. These systems decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta transitions. Only <sup>79</sup>Ni and <sup>80</sup>Ni have measured half-lives. The half-life values of the <sup>79</sup>Ni – <sup>80</sup>Ni systems agree with the data [39]. Model calculations also predict the correct decay mode of these  $2d_{5/2}$  nickel nuclei.

#### 95 ≥ A ≥ 81 Nickel Isotopes without Experimental Half-Life Data

The  $a(A)$  values for  $81 \geq A \geq 95$  nickel isotopes were obtained from a linear fit based on the half-lives of <sup>79</sup>Ni – <sup>80</sup>Ni. The resulting  $a(A)$  values are listed in Table 1.

<sup>81</sup>Ni – <sup>84</sup>Ni complete the  $2d_{5/2}$  neutron shell, and decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transitions. These nickel systems have beta decay half-lives between 4.25 and 10.5 ms. Single-particle calculations are within about 25% of the Japanese Data Compilation calculations [39].

<sup>85</sup>Ni – <sup>92</sup>Ni nuclei fill the  $1g_{7/2}$  neutron shell, and have beta decay half-lives in the range of 0.538 to 3.06 ms. These systems decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transitions. The model results for <sup>85</sup>Ni and <sup>86</sup>Ni are within about 35% of the Japanese Data Compilation calculations [39]. <sup>86</sup>Ni is the last nickel system included in the Japanese Data Compilation calculations [39].

<sup>93</sup>Ni and <sup>94</sup>Ni fill the  $3s_{1/2}$  neutron shell, and these systems decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transitions. The half-lives of <sup>93</sup>Ni and <sup>94</sup>Ni are 0.444 and 0.369 ms, respectively. The  $2d_{3/2}$  neutron shell is partially filled by the <sup>95</sup>Ni system. This system decays through an allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transition, and has a half-life of 0.310 ms.

No nickel systems with  $A > 95$  are predicted by either the model or the Japanese Data Compilation calculations [39]. This model limitation occurs because only 67 neutrons are bound in nickel

system.

## Conclusions

Within to are the scope of the proposed single particle model, neutron excess nickel isotopes terminate with <sup>95</sup>Ni. The model predicts that the  $81 \leq A \leq 95$  nickel systems have beta decay half-lives in the range of 0.310 – 10.5 ms. These neutron excess nickel systems decay through allowed  $2p_{1/2}(n)$  to  $2p_{3/2}(p)$  beta decay transitions. The model likely overestimates the actual half-life values, because it does not include the short-lived neutron emission decay modes.

## References

1. Siegel, D. M., & Metzger, B. D. (2017). Three-dimensional general-relativistic magnetohydrodynamic simulations of remnant accretion disks from neutron star mergers: Outflows and r-process nucleosynthesis. *Phys Rev Lett*, 119, 231102.
2. National Academy of Sciences. (2007). Scientific opportunities with a rare isotope facility in the United States (Report No. 11796). National Research Council.
3. Fukuda, N., et al. (2018). Identification of new neutron-rich isotopes in the rare-earth region produced by 345 MeV/nucleon <sup>238</sup>U. *Journal of the Physical Society of Japan*, 87, 014202.
4. Shimizu, Y., et al. (2018). Observation of new neutron rich isotopes among fission fragments from in flight fission of 345 MeV/nucleon <sup>238</sup>U: Search for new isotopes conducted concurrently with decay measurement campaigns. *Journal of the Physical Society of Japan*, 87, 014203.
5. Kurcewicz, J., et al. (2012). Discovery and cross section measurement of neutron rich isotopes in the element range from neodymium to platinum with the FRS. *Physics Letters B*, 717, 371–376.
6. Baumann, T., et al. (2007). Discovery of <sup>40</sup>Mg and <sup>42</sup>Al suggests neutron drip line slant towards heavier isotopes. *Nature*, 449, 1022–1024.
7. Tarasov, O. B., et al. (2013). Production cross sections from <sup>82</sup>Se fragmentation as indications of shell effects in neutron rich isotopes close to the drip line. *Physical Review C*, 87, 054612.
8. Bevelacqua, J. J. (2018). Decay characteristics of neutron excess calcium nuclei. *Physics Essays*, 31(4), 462–469.
9. Bevelacqua, J. J. (2020). Decay characteristics of neutron excess iron nuclei. *Physics Essays*, 32(2), 175–181.
10. Bevelacqua, J. J. (2020). Decay characteristics of neutron excess fluorine nuclei (24XLL9). *Qeios*, Article 1. <https://doi.org/10.32388/24XLL9>
11. Bevelacqua, J. J. (2020). Decay characteristics of neutron

- excess zinc nuclei (JZI1LG). Qeios, Article 1. <https://doi.org/10.32388/JZI1LG>
12. Bevelacqua, J. J. (2021). Decay characteristics of neutron excess neon nuclei (1WR291). Qeios, Article 1. <https://doi.org/10.32388/1WR291>
  13. Bevelacqua, J. J. (2021). Decay characteristics of neutron excess sodium nuclei (1Y819A). Qeios, Article 1. <https://doi.org/10.32388/1Y819A>
  14. Bevelacqua, J. J. (2021). Decay characteristics of neutron excess magnesium nuclei (KIB58L). Qeios, Article 1. <https://doi.org/10.32388/KIB58L>
  15. Bevelacqua, J. J. (2022). Decay characteristics of neutron excess aluminum nuclei (LCAO3W). Qeios, Article 1. <https://doi.org/10.32388/LCAO3W>
  16. Bevelacqua, J. J. (2022). Decay characteristics of neutron excess silicon nuclei (Y6HDZF). Qeios, Article 1. <https://doi.org/10.32388/Y6HDZF>
  17. Bevelacqua, J. J. (2023). Decay characteristics of neutron excess phosphorus nuclei (Z16MGO). Qeios, Article 1. <https://doi.org/10.32388/Z16MGO>
  18. Bevelacqua, J. J. (2023). Decay characteristics of neutron excess sulfur nuclei (QO9K3E). Qeios, Article 1. <https://doi.org/10.32388/QO9K3E>
  19. Bevelacqua, J. J. (2023). Decay characteristics of neutron excess chlorine nuclei (HXV1XN). Qeios, Article 1. <https://doi.org/10.32388/HXV1XN>
  20. Bevelacqua, J. J. (2023). Decay characteristics of neutron excess argon nuclei (JDLHDL). Qeios, Article 1. <https://doi.org/10.32388/JDLHDL>
  21. Bevelacqua, J. J. (2024). Decay characteristics of neutron excess potassium nuclei (RBF GK2). Qeios, Article 1. <https://doi.org/10.32388/RBF GK2>
  22. Bevelacqua, J. J. (2024). Decay characteristics of neutron excess scandium nuclei (25NGQR). Qeios, Article 1. <https://doi.org/10.32388/25NGQR>
  23. Bevelacqua, J. J. (2024). Decay characteristics of neutron excess titanium nuclei (NFSVCP). Qeios, Article 1. <https://doi.org/10.32388/NFSVCP>
  24. Bevelacqua, J. J. (2024). Decay characteristics of neutron excess vanadium nuclei (9VY02M). Qeios, Article 1. <https://doi.org/10.32388/9VY02M>
  25. Bevelacqua, J. J. (2024). Decay characteristics of neutron excess chromium nuclei (K7DJZP). Qeios, Article 1. <https://doi.org/10.32388/K7DJZP>
  26. M. Terasawa, K. Sumiyosh, T. Kajino, G. J. Mathews, and I. Tanihata, New Nuclear Reaction Flow during r-Process Nucleosynthesis in Supernovae: Critical Role of Light Neutron-Rich Nuclei (2001). THE ASTROPHYSICAL JOURNAL, 562, 470.
  27. Lukasiak, A., & Sobieczewski, A. (1975). Estimations of half lives of far superheavy nuclei with  $Z \approx 154-164$ . Acta Physica Polonica B, 6, 147–155.
  28. Petrovich, F., Philpott, R. J., Robson, D., Bevelacqua, J. J., Golin, M., & Stanley, D. (1976). Comments on primordial superheavy elements. Physical Review Letters, 37, 558–561.
  29. Bevelacqua, J. J. (2025). Decay characteristics of neutron excess manganese nuclei (FKQ489). Qeios, Article 1. <https://doi.org/10.32388/FKQ489>
  30. Bevelacqua, J. J. (2025). Decay characteristics of neutron excess cobalt nuclei. Sci Set Journal of Physics, 4(3), 1–4.
  31. Brown, G. E., Gunn, J. H., & Gould, P. (1963). Effective mass in nuclei. Nuclear Physics, 46, 598–604.
  32. Fox, L., & Godwin, E. T. (1949). Some new methods for the numerical integration of ordinary differential equations. Proceedings of the Cambridge Philosophical Society, 45, 373–385.
  33. Rost, E. (1968). Proton shell model potentials for lead and the stability of superheavy nuclei. Physics Letters, 26B, 184–190.
  34. Blomqvist, J., & Wahlborn, S. (1959). Shell model calculations in the lead region with a diffuse nuclear potential. Arkiv för Fysik, 16, 545–558.
  35. Ray, L., & Hodgson, P. E. (1979). Neutron densities and the single particle structure of several even even nuclei from Ca40 to Pb208. Physical Review C, 20, 2403–2415.
  36. Schwierz, N., Wiedenhöfer, I., & Volya, A. (2007). Parameterization of the Woods Saxon potential for shell model calculations. arXiv preprint arXiv:0709.3525 [nucl th].
  37. Baum, E. M., Ernesti, M. C., Knox, H. D., Miller, T. R., & Watson, A. M. (2010). Nuclides and Isotopes – Chart of the Nuclides (17th ed.). Knolls Atomic Power Laboratory.
  38. National Nuclear Data Center. (2024). NuDat3 (Nuclear Structure and Decay Data). Brookhaven National Laboratory.
  39. Koura, H., et al. (2018). Chart of the Nuclides 2018. Japanese Nuclear Data Committee & Nuclear Data Center, Japan Atomic Energy Agency.
  40. Wong, C. Y. (1966). Additional evidence of stability of the superheavy element  $^{310}126$  according to the shell model.