

# Euler-Lagrange Equations on Three-Dimensional Almost Kenmotsu Manifolds

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## Abstract

In this study we concluded the Euler- Lagrange equations on  $(M^3, \phi, \xi, \eta, g)$ , The Three-Dimensional almost Kenmotsu manifolds, have been derived also important applications of Euler-Lagrange mechanical systems. Finally achieved that Three-Dimensional almost Kenmotsu manifolds have this system in Mechanics and Physical Fields as well as in differential geometry.

**Keywords:** Kenmotsu Manifolds, Three – Dimensional Almost Kenmotsu Manifolds, Euler-Lagrange Equations

## Introduction

The geometric study of dynamical systems is an important chapter of contemporary mathematics due to its applications in Mechanics, Theoretical Physics.

There are also a large number of studies on this subject, for example Jun, J and Pathak submitted On Kenmotsu manifolds [1]. and De, U.C. and Pathak obtained on 3-dimensional Kenmotsu manifolds [2].

It is some important work for examples [3-7]. In this paper we will study Kenmotsu manifolds

In this paper, we Euler-Lagrange Equations on Three-Dimensional almost Kenmotsu manifolds. After Introduction in Section 1, we consider Historical Background paper basic. Section 2 deals with the study preliminary. Section 3 is devoted to study 3. three – dimensional Almost Kenmotsu manifolds. Section 4 is devoted to study . Euler-Lagrange Equations on Three-Dimensional almost Kenmotsu manifolds

## Preliminary

In this in this preliminary chapter, we recall basic definitions, results and formulas which we shall use in the subsequent chapters of the paper.

## Definition (Kenmotsu Manifolds)

Let  $M^{(2n+1)}$  be a  $(2n+1)$ -dimensional smooth differentiable manifold  $(\phi, \xi, \eta, g)$  be an almost contact Riemannian manifold. where  $\phi$  is a  $(1-1)$  tensor field  $\eta$  is a 1- form and the Riemannian metric .It well known that

$$\begin{aligned} \phi(\xi) &= 0 \\ \eta(\phi(x)) &= 0 \quad \text{and} \quad \eta(\xi) = 1 \\ \phi^2(X) &= -X + \eta(X)\xi, \quad \phi^2 = -1 + \eta \otimes \xi \\ g(X, \xi) &= \eta(X) \\ g(\phi x, \phi y) &= g(x, y) - \eta(x)\eta(y) \\ \text{rank } \phi &= n - 1 \end{aligned}$$

The fundamental 2- form of an almost contact metric manifolds is defined by

$$\phi(x, y) = g(x, \phi y)$$

Lemma 2.2 [2] If  $\omega, \theta$  and  $k$ -form be respectively then

- $d\omega \wedge \theta = -d\theta \wedge \omega$
- $d = d \circ d = 0$
- $d(\omega \wedge \psi) = d\omega \wedge \psi + (-1)^k \omega \wedge d\psi$

Lemma 2.3 [2] Suppose  $(U, x_1, \dots, x_n)$  is a chart on a manifold. Then

$$\left( \frac{\partial x^j}{\partial x^i} \right) = \delta_j^i = \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases}$$

### Three – Dimensional Almost Kenmotsu Manifolds

**Definition 3.1** Let.  $(M^3, \phi, \xi, \eta, g)$  be a 3-dimensional almost Kenmotsu manifold. In what follows, we denote by  $L=R(\cdot, \xi)\xi$ ,  $h=1/2 L\xi\phi$  and  $h^\perp=h\phi$ , where  $L$  denotes the Lie differentiation and  $R$  is the Riemannian curvature tensor. From Dileo and Pastore [6, 8], we see that both  $h$  and  $h^\perp$  are symmetric operators and we recall some properties of almost Kenmotsu manifolds as follows:

$$h\xi = L\xi = 0, \quad \text{tr } h = \text{tr}(h^\perp) = 0$$

$$h^\perp \phi = \phi h = 0$$

$$\nabla \xi = h + id - \eta \circ \phi$$

**Proposition 3.2** Any 3-dimensional almost Kenmotsu manifold is Kenmotsu if and only if  $h$  vanishes.

#### Definition 3.3

Let three-dimensional manifold  $M=f(x,y,z)$   $R^3, z \neq 0$ ; where  $(x,y,z)$  are the standard coordinates in  $R^3$ : The vector fields

$$e_1 = \frac{\partial}{\partial x}, \quad e_2 = \frac{\partial}{\partial y}, \quad e_3 = \frac{\partial}{\partial z}$$

are linearly independent at each point of  $M$ : Let  $g$  be the Riemannian metric defined by

$$g(e_1, e_3) = g(e_2, e_3) = g(e_1, e_2) = 0 \\ g(e_1, e_1) = g(e_2, e_2) = g(e_3, e_3) = 1$$

Let  $\eta$  be the 1-form defined by  $\eta(Z) = g(Z, e_3)$  for any  $Z \in \mathfrak{X}(M)$ .

#### Proposition 3.4

Let  $\phi$  be the  $(1,1)$  tensor field defined by

$$\phi(e_1) = -e_2 \\ \phi(e_2) = e_1 \\ \phi(e_3) = 0$$

Then using the linearity of  $\phi$  and  $g$  we have

$$\eta(e_3) = 1; \quad \phi^2(Z) = -Z + \eta(Z)e_3; \\ g(\phi Z, \phi W) = g(Z, W) - \eta(Z)\eta(W);$$

for any  $Z, W \in \mathfrak{X}(M)$ : Thus for  $e_3 = \xi$ ,  $(\phi, \xi, \eta, g)$  defines an almost contact metric structure on  $M$ : Now, by direct computations we obtain

$$[e_1, e_3] = 0; \quad [e_2, e_3] = -e_2, \quad [e_1, e_3] = -e_1$$

#### Proposition 3.5

The vector fields

$$e_1 = \frac{\partial}{\partial x}, \quad e_2 = \frac{\partial}{\partial y}, \quad e_3 = \frac{\partial}{\partial z}$$

If  $\phi$  is defined a complex manifold  $M$  then  $\phi^2 = \phi \circ \phi = -I$   
Proof.

$$\phi^2(e_1) = \phi(\phi(e_1)) = \phi(-e_2) = -\phi(e_2) = -e_1$$

$$\phi^2(e_2) = \phi(\phi(e_2)) = \phi(e_1) = \phi(e_1) = -e_2$$

As can  $\phi^2$  is  $-I$  (complex) or 0

### Euler-Lagrange Equations on Three-Dimensional Almost Kenmotsu Manifolds

In this section, we shall obtain the version Euler-Lagrange equations for classical mechanics structured with Three-Dimensional almost Kenmotsu manifolds introduced in Let semispray be a vector field as follows

$$\xi = x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + z \frac{\partial}{\partial z}, \quad x = \dot{x}, \quad y = \dot{y}, \quad z = \dot{z}$$

By Liouville vector field on Three-Dimensional almost Kenmotsu manifolds space form  $(M^3, \phi, \xi, \eta, g)$ , we call the vector field determined by  $V = \phi \xi$  and calculated by

$$\phi \xi = \phi \left( x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + z \frac{\partial}{\partial z} \right) = x \phi \left( \frac{\partial}{\partial x} \right) + y \phi \left( \frac{\partial}{\partial y} \right) + z \phi \left( \frac{\partial}{\partial z} \right)$$

$$\phi \xi = -x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x} + z(0) = -x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x} + z(0)$$

$$\phi \xi = -x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x}$$

Denote  $T$  by the kinetic energy and  $P$  by the potential energy of mechanics system on Three-Dimensional almost Kenmotsu manifolds. Then we write by  $L = T - P$  Lagrangian function and by  $E_L = V(L) - L$

the energy function associated  $L$ . Operator  $i_{\phi} L$  defined by  $i_{\phi} L = L(\phi) - L$

is called the interior product with  $\phi$ , or sometimes the insertion operator, or contraction by  $\phi$ . The exterior vertical derivation  $d_{\phi}$  is defined by

$$d_{\phi} = [i_{\phi}, d] = i_{\phi} d - di_{\phi}$$

where  $d$  is the usual exterior derivation? For almost product structure  $\phi$  determined by the closed Three-Dimensional almost Kenmotsu manifolds form is the closed 2-form given by  $\phi_L = -dd_{\phi} L$  such that

$$d_{\phi} = -x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x}$$

$$d_{\phi} L = \left( -x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x} \right) L = -x \frac{\partial L}{\partial y} + y \frac{\partial L}{\partial x}$$

Thus we get

$$\phi_L = -d(d_{\phi} L) = -d \left( -x \frac{\partial L}{\partial y} + y \frac{\partial L}{\partial x} \right)$$

$$\phi_L = X \frac{\partial^2 L}{\partial x \partial y} dx \wedge dx + X \frac{\partial^2 L}{\partial y \partial y} dy \wedge dy - Y \frac{\partial^2 L}{\partial x \partial x} dx \wedge dy - Y \frac{\partial^2 L}{\partial y \partial y} dy \wedge dy - Z \frac{\partial^2 L}{\partial x \partial z} dx \wedge dz \\ - Z \frac{\partial^2 L}{\partial y \partial z} dy \wedge dz$$

Because of the closed Three-Dimensional almost Kenmotsu manifolds form  $\phi_L$  on Three-Dimensional almost Kenmotsu manifolds space form  $(M^3, \phi, \xi, \eta, g)$  is para-symplectic structure, one may obtain

$$E_L = X \frac{\partial L}{\partial Y} + Y \frac{\partial L}{\partial X} - L$$

Considering (0,1) we calculate

$$dE_L = d\left(X\frac{\partial L}{\partial Y} + Y\frac{\partial L}{\partial X} - L\right)$$

$$dE_L = X\frac{\partial^2 L}{\partial x \partial y} dx + y\frac{\partial^2 L}{\partial y \partial y} dx + X\frac{\partial^2 L}{\partial x \partial x} dy + Y\frac{\partial^2 L}{\partial y \partial x} dy - \frac{\partial L}{\partial x} dx - Z\frac{\partial^2 L}{\partial x \partial z} dx - Z\frac{\partial^2 L}{\partial y \partial z} dy - \frac{\partial L}{\partial y} dy$$

Taking care of  $i_{\xi} \phi^* L = dE_L$ , we have

$$X\frac{\partial^2 L}{\partial x \partial y} dx + y\frac{\partial^2 L}{\partial y \partial y} dx + Z\frac{\partial^2 L}{\partial y \partial z} dx - \frac{\partial L}{\partial x} dx + X\frac{\partial^2 L}{\partial x \partial x} dy + Y\frac{\partial^2 L}{\partial y \partial x} dy + Z\frac{\partial^2 L}{\partial x \partial z} dy + \frac{\partial L}{\partial y} dy = 0$$

$$\frac{\partial L}{\partial y} \left( X\frac{\partial}{\partial x} dx + y\frac{\partial}{\partial y} dx + Z\frac{\partial}{\partial z} dx \right) - \frac{\partial L}{\partial x} dx + \frac{\partial L}{\partial x} \left( X\frac{\partial}{\partial x} dy + Y\frac{\partial}{\partial y} dy + Z\frac{\partial}{\partial z} dy \right) + \frac{\partial L}{\partial y} dy = 0$$

$$\frac{\partial L}{\partial y} \left( X\frac{\partial}{\partial x} + y\frac{\partial}{\partial y} + Z\frac{\partial}{\partial z} \right) dx - \frac{\partial L}{\partial x} dx + \frac{\partial L}{\partial x} \left( X\frac{\partial}{\partial x} + Y\frac{\partial}{\partial y} + Z\frac{\partial}{\partial z} \right) dy + \frac{\partial L}{\partial y} dy = 0$$

$$\left[ \frac{\partial L}{\partial y} \left( X\frac{\partial}{\partial x} + y\frac{\partial}{\partial y} + Z\frac{\partial}{\partial z} \right) - \frac{\partial L}{\partial x} \right] dx + \left[ \frac{\partial L}{\partial x} \left( X\frac{\partial}{\partial x} + Y\frac{\partial}{\partial y} + Z\frac{\partial}{\partial z} \right) + \frac{\partial L}{\partial y} \right] dy = 0$$

If the curve  $\alpha: I \rightarrow M^3$  be integral curve of  $\xi$ ,

$$\alpha = \frac{\partial}{\partial t} = X\frac{\partial}{\partial x} + y\frac{\partial}{\partial y} + Z\frac{\partial}{\partial z}$$

which satisfies

$$\left[ \frac{\partial}{\partial t} \frac{\partial L}{\partial y} - \frac{\partial L}{\partial x} \right] dx + \left[ \frac{\partial}{\partial t} \frac{\partial L}{\partial x} + \frac{\partial L}{\partial y} \right] dy = 0$$

it follows equations

$$\frac{\partial}{\partial t} \frac{\partial L}{\partial y} - \frac{\partial L}{\partial x} = 0, \quad \frac{\partial}{\partial t} \frac{\partial L}{\partial x} + \frac{\partial L}{\partial y} = 0$$

so-called Euler-Lagrange equations whose solutions are the paths of the semispray  $\xi$  on Three-Dimensional almost Kenmotsu manifolds space form  $(M^3, \phi, \xi, \eta, g)$ . Finally, one may say that the triple  $(M^3, \phi, \xi, \eta, g)$  is mechanical system on Three-Dimensional almost Kenmotsu manifolds  $(M^3, \phi, \xi, \eta, g)$ . Therefore we say

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